

Available online at www.sciencedirect.com



Journal of Molecular Spectroscopy 237 (2006) 149-162

Journal of MOLECULAR SPECTROSCOPY

www.elsevier.com/locate/jms

# High-resolution Fourier-transform spectroscopy of <sup>18</sup>O-enriched water molecule in the 1080–7800 cm<sup>-1</sup> region

An-Wen Liu, Jun-He Du, Ke-Feng Song, Le Wang, Lei Wan, Shui-Ming Hu \*

Hefei National Laboratory for Physical Sciences at Microscale, University of Science and Technology of China, Hefei 230026, China Shanghai Institute for Advanced Studies, University of Science and Technology of China, Shanghai 201315, China

> Received 24 February 2006; in revised form 10 March 2006 Available online 28 March 2006

#### Abstract

Fourier-transform absorption spectrum of <sup>18</sup>O enriched water sample was recorded in the 1080–7800 cm<sup>-1</sup> region. The transitions of H<sub>2</sub><sup>18</sup>O were assigned on the base of the high accuracy ab initio calculations by Partridge and Schwenke (PS). One thousand two hundred and forty-six ro-vibrational energy levels were retrieved, which belong to the  $v_1 + \frac{1}{2}v_2 + v_3 = 0.5$ , 1, and 1.5 polyads: (010), (100), (020), (001), (110), (011), and (030) states. Four hundred and thirty-two of them are reported for the first time. The results are also investigated comparing with the PS calculations and the available literature data in the considered spectral range. HD<sup>18</sup>O as a small concentration in the sample, more than 4900 transitions were also observed. These transitions yield 1066 ro-vibrational energies of nine vibrational states, 504 energy levels were obtained for the first time.

© 2006 Elsevier Inc. All rights reserved.

Keywords: Vibration-rotation spectroscopy; Water vapor absorption; Spectroscopic parameters

## 1. Introduction

As the second primary isotope of the water molecule,  $H_2^{18}O$  and its infrared absorption contributes noticeably to the atmospheric absorption and might be taken into account in the explanation of atmospheric radiation balance. In the last few decades, numerous spectroscopic studies on the H<sub>2</sub><sup>18</sup>O molecule have been carried out with microwave [1-5], Fourier-transform infrared [6-24], and cavity ring-down spectroscopy [25,26] in the near-IR and visible region. More than 800 energy levels have been reported from our previous Fourier-transform study [27,28]. All these results including  $H_2^{18}O$  energy levels, transition frequencies, and strengths are available in the most complete database to date [29]. Despite of all the above undertakings, the high resolution study of  $H_2^{18}O$ below 1 µm region is not yet complete, compared with the detailed simulation given by Partridge and Schwenke

\* Corresponding author. Fax: +86 551 3602969.

E-mail address: smhu@ustc.edu.cn (S.-M. Hu).

(PS) [30,31]. Many lines predicted with relatively strong and middle intensities are not yet observed. For the HD<sup>18</sup>O molecule, to the best of our knowledge, only a few bands including  $v_1$ ,  $v_2$ ,  $2v_2$ ,  $v_3$ ,  $3v_2$ , and  $v_1 + v_2^{-1}$  have been studied [32–34].

The purpose of this work is, besides of a continuation to our previous studies of the high resolution Fourier-transform absorption spectroscopy of  $H_2^{18}O$  [27,28], the study of HD<sup>18</sup>O in the region between 1080 and 7800 cm<sup>-1</sup> will be also presented. The experimental details, transition assignments, and analysis will be given in the following sections.

#### 2. Experimental details

The <sup>18</sup>O enriched water sample was purchased from Aldrich Chemical Company, Inc. The stated isotopic concentration of the <sup>18</sup>O atom is 95%. The absorption

<sup>0022-2852/\$ -</sup> see front matter @ 2006 Elsevier Inc. All rights reserved. doi:10.1016/j.jms.2006.03.008

 $<sup>1</sup> v_1$  and  $v_3$  present the OD and OH stretching modes of HDO, respectively.  $v_2$  is for the bending mode.

spectra were recorded at room temperature with a Bruker IFS 120HR Fourier-transform spectrometer equipped with a path length adjustable multi-pass gas cell. Because of the wide spectral range and the large variation of the absorption line intensities, different experimental conditions were used in the measurements as listed in Table 1. The pressure was measured using two capacitance manometers of 200 Pa and 133 hPa full-scale range with an overall accuracy of 0.5%. In most measurements, optical filters were applied to increase the signal–noiseratio and to allow the high spectral resolution measurements. The line positions were calibrated with water lines given by the HITRAN2004 [29] database. The accuracy of the line positions of unblended and not-very-weak lines was estimated to be better than  $0.0012 \text{ cm}^{-1}$ . Parts of two spectrum together with the assignments are presented in Figs. 1 and 2.

Because we could not substitute all the water adsorbed on the walls of the sample cell with the very limited amount of sample, it is difficult to accurately

Table 1

Experimental conditions in the Fourier-transform measurements of <sup>18</sup>O enriched water sample

Region (cm <sup>-1</sup> )	Source	Beam splitter	Detector	Path length (m)	Pressure (Pa)	Resolution (cm <sup>-1</sup> )
950-1350	Globar	KBr	MCT	15	1307	0.005
1200-1700	Globar	KBr	MCT	15	126	0.005
1200-1700	Globar	KBr	MCT	15	1307	0.005
800-5000	Globar	KBr	MCT	15	1307	0.008
2500-3350	Tungsten	$CaF_2$	InSb	15	126	0.005
2500-3350	Tungsten	$CaF_2$	InSb	15	1307	0.005
2500-3350	Tungsten	$CaF_2$	InSb	105	1307	0.008
3200-4300	Tungsten	$CaF_2$	InSb	15	126	0.008
3200-4300	Tungsten	$CaF_2$	InSb	15	1307	0.008
3900-5100	Tungsten	$CaF_2$	InSb	105	1307	0.01
4900-6400	Tungsten	$CaF_2$	InSb	15	11	0.01
4900-6400	Tungsten	$CaF_2$	InSb	15	126	0.01
4900-6400	Tungsten	$CaF_2$	InSb	15	1307	0.01
4900-6400	Tungsten	$CaF_2$	InSb	105	1307	0.01
5000-11000	Tungsten	$CaF_2$	Ge	15	215	0.015
5000-11000	Tungsten	$CaF_2$	Ge	15	1307	0.015
5000-11000	Tungsten	CaF <sub>2</sub>	Ge	105	1307	0.015



Fig. 1. A small part of the  $H_2^{18}O$  spectrum and the line assignments of  $v_1$ ,  $v_3$ ,  $2v_2$ , and  $3v_2$  bands marked by square, triangle, circle, and diamond, respectively. Experimental conditions: absorption path length 105 m, sample pressure 1307 Pa, unapodized resolution 0.008 cm<sup>-1</sup>, and 0.01 cm<sup>-1</sup> for upper and lower panels, respectively. Lines marked with <sup>16</sup>O and <sup>17</sup>O are due to  $H_2^{16}O$  and  $H_2^{17}O$ .



Fig. 2. A small part of the HD<sup>18</sup>O spectrum. (Upper panel) 5170–5175 cm<sup>-1</sup> region of the  $v_2 + v_3$  band recorded with InSb detector and 0.01 cm<sup>-1</sup> unapodized resolution. (Lower panel) 7060–7065 cm<sup>-1</sup> region of the  $2v_3$  band recorded at 0.015 cm<sup>-1</sup> unapodized resolution. Experimental conditions: absorption path length 105 m, sample pressure 1307 Pa. Lines marked with <sup>16</sup>O, <sup>18</sup>O, <sup>17</sup>O, and D are due to H<sub>2</sub><sup>16</sup>O, H<sub>2</sub><sup>18</sup>O, H<sub>2</sub><sup>17</sup>O, and HD<sup>16</sup>O, respectively.



Fig. 3. H<sub>2</sub><sup>18</sup>O stick spectrum from HITRAN2004 database, Schwenke and Partridge calculations [30,31] and this work.

determine the real isotopic concentration of the sample used in our measurements. As discussed in [27], the isotope abundance of  $H_2^{18}O$  in each spectrum were estimated by the comparison of the line intensities (retrieved from line profile fittings to our spectrum) of some "moderate" lines with their values given in the HITRAN2004 database. Then this isotopic abundance value from each spectrum will be used to determine the intensities of other lines. The intensities in this work were derived as in [27]. The precision of the intensities of well isolated and not-very-weak  $H_2^{18}O$  lines is estimated to be about 15% which should mainly comes from the ambiguity of the partial pressure of the sample.

## 3. Spectral analysis and results

$$3.1. H_2^{18}O$$

With the help of the high accurate calculations by Schwenke and Partridge [30,31], the observed lines of  $H_2^{18}O$  in the first and second hexade region [27,28] have been successfully assigned. Seven bands of H<sub>2</sub><sup>18</sup>O locate in the 1080–6000 cm<sup>-1</sup> region:  $v_2$ ,  $v_1$ ,  $2v_2$ ,  $v_3$ ,  $v_1 + v_2$ ,  $3v_2$ , and  $v_2 + v_3$ . The assignments of the H<sub>2</sub><sup>18</sup>O lines in this region were done straightforwardly by the identification procedure in a similar way. The final H<sub>2</sub><sup>18</sup>O identification list is attached to the paper as the Supplementary Material I, which includes 5031 lines corresponding to 5234 transitions. For each line, the experimental and PS calculated intensities are given followed by ro-vibrational assignments. It also contains some  $H_2^{18}O$  lines superimposed by the transitions of other water isotopes, these lines are specially marked in the list.

In a new version of HITRAN2004, there are 4297 transitions stronger than  $2.5 \times 10^{-5} \text{cm}^{-1}$  /(atm cm) assigned to  $\text{H}_2^{18}\text{O}$  in the studied region. The detection limit of the spectrum presented in this work extends to about  $10^{-6} \text{ cm}^{-1}$  /(atm cm). On the one hand, there are 1300 transitions included in HITRAN2004 database but

Table 2	
Summary of the determined	energy levels of H <sub>2</sub> <sup>18</sup> O

V	$J^{\max}$	$K_a^{\max}$	Number of levels					
			This work	[34]	New	Diff		
010	20	10	205	192	32	26		
020	16	10	148	98	50	2		
100	15	10	197	126	71	3		
001	19	11	223	161	62	3		
030	15	8	117	3	114	0		
110	15	8	147	102	46	7		
011	18	10	209	153	57	10		
021	11	7	75	90				
120	9	4	18	59				
Total			1339	984	432	51		

Table 3 New ro-vibrational energy levels of  $H_2^{18}O$  (cm<sup>-1</sup>)

		0,		2 ( )		
V	J	$K_a$	$K_c$	Ε	δ	Ν
010	11	10	2	4730.2169		1
010	11	10	1	4730.2171		1
010	12	8	4	4528.0225		1
010	12	9	4	4769.5012	0.7	2
010	12	9	3	4769.5012	0.7	2
010	12	10	3	5024.7743		1
010	12	10	2	5024.7743		1
010	13	6	8	4416.2162	0.8	3
010	13	7	7	4618.0661		1
010	13	7	6	4618.6401	0.6	3
010	13	8	6	4842.0717		1
010	13	8	5	4842.1121		1
010	14	4	10	4511.3382		1
010	14	5	9	4623.2336		1
010	14	6	9	4/52.4/38		1
010	14	6	8	4/64./951		1
010	14	/	12	4934.0021	0.1	1
010	15	4	12	4/1/.4/24	0.1	2
010	15	4	11	4877.4065		1
010	16	2	14	4827.4905		1
010	16	3	14	5069 8538		1
010	16	4	13	5073 9073		1
010	17	1	16	4882 2282		1
010	17	2	16	4882 2539		1
010	17	2	15	5183 4974		1
010	18	1	17	5236 7778		1
010	18	2	17	5236.7759		1
010	19	0	19	5220.0202		1
010	19	1	19	5220.0185		1
010	19	1	18	5608.6196		1
010	20	1	20	5588.0767		1
020	7	7	1	4780.7990	0.4	4
020	7	7	0	4780.7989	0.5	5
020	8	7	2	4977.3155	0.7	5
020	8	7	1	4977.3155	0.9	3
020	8	8	1	5225.4954		1
020	8	8	0	5225.4955		1
020	9	4	6	4579.8581	0.9	5
020	9	5	5	4759.7437	0.5	6
020	9	6	4	4967.2336	0.6	3
020	9	6	3	4966.5170	0.7	5
020	9	7	3	5197.6542		1
020	9	7	2	5197.6747	0.4	2
020	9	8	2	5447.0708		1
020	9	8	1	5447.0702		1
020	9	9	1	5709.4806		1
020	9	9	0	5/09.4806	0.2	1
020	10	2	0 7	4020.9021	0.2	2
020	10	3	7	4735.0014	0.5	2
020	10	4	6	4844 7631	0.5	1
020	10	- <del>1</del> 5	6	5002 9269	0.6	3
020	10	5	5	5005 8557	0.5	4
020	10	6	5	5210 5721	0.5	- <del>-</del>
020	10	6	4	5209 9748		1
020	10	7	4	5441.5891		1
020	10	, 9	2	5956.1109		1
020	11	2	10	4697.2244	0.6	4
020	11	2	9	4887.0990	0.4	5
020	11	3	9	4902.0883	0.9	3

*Note.* N is the number of lines used for the upper level determination and  $\delta$  denotes the experimental uncertainties (combination differences) in  $10^{-3}$  cm<sup>-1</sup>.

Table 3 (continued)

Table 3	(continued	<i>!</i> )					Table 3	(continued	l)				
V	J	$K_a$	$K_c$	Ε	δ	Ν	V	J	$K_a$	$K_c$	Ε	δ	Ν
020	11	3	8	5016.8388		1	100	13	5	8	6226.8674		1
020	11	4	8	5084.1590	0.2	2	100	13	6	8	6341.5623		1
020	11	4	7	5124.9905		1	100	13	6	7	6345.4345		1
020	11	5	7	5269.4179		1	100	13	7	7	6506.2070		1
020	11	5	6	5276.3759	0.4	2	100	13	7	6	6506.9693		1
020	12	1	11	4946.5352	0.5	4	100	13	8	6	6697.3398		1
020	12	2	10	5162.8585	0.0	2	100	14	0	14	5674.6693		1
020	12	4	9	5367.7776		1	100	14	1	14	5674.6700	0.2	2
020	12	4	8	5431.5156		1	100	14	1	13	5923.3917		1
020	12	5	8	5557.6643		1	100	14	2	13	5923.5386		1
020	13	1	12	5214.6887	0.3	3	100	14	2	12	6143.30//		1
020	13	2	12	5215.1013		1	100	14	2	12	6220 2054		1
020	13	2	11	5455.5986		1	100	14	3	11	6333 0412		1
020	13	3	10	5461.0621		1	100	14	4	10	6502 0865		1
020	15	3	10	5182 0217	0.6	1	100	14	5	9	6575 1397		1
020	14	1	14	5182.9217	0.0	2	100	14	6	9	6656 6080	0.3	2
020	14	3	14	5767 1183	0.5	1	100	14	6	8	6685 8146	0.5	1
020	14	3	12	6322 1378		1	100	14	7	8	6838 7827		1
020	15	0	12	5758 0375		1	100	14	, 7	7	6840 6532		1
020	16	1	16	5758 0380		1	100	15	0	15	5952.8898	0.7	2
100	9	8	2	5595 7977	03	2	100	15	1	15	5952.8881	017	- 1
100	9	8	1	5595 7975	0.5	2	100	15	2	13	6455.7361		1
100	9	9	1	5803 1037	1.0	2	100	15	3	13	6456.4584		1
100	9	9	0	5803.1038	0.9	2	100	15	3	12	6663.1368		1
100	10	6	4	5476.8202	0.9	4	100	15	4	11	6829.7715		1
100	10	7	3	5646.2254		1	100	15	7	8	7198.1039		1
100	10	8	3	5836.6874		1	001	9	9	1	5858.8783	0.9	3
100	10	8	2	5836.6890	0.3	2	001	9	9	0	5858.8777	0.6	3
100	10	9	2	6045.6389	0.9	2	001	10	9	2	6101.8094	1.0	3
100	10	9	1	6045.6393	0.8	2	001	10	9	1	6101.8093	0.6	3
100	10	10	1	6276.3779		1	001	10	10	1	6313.4094		1
100	10	10	0	6276.3757		1	001	10	10	0	6313.4094		1
100	11	3	9	5306.6933	0.2	4	001	11	7	5	5984.0698	0.1	2
100	11	4	8	5450.1317	0.5	2	001	11	7	4	5984.2417	0.4	3
100	11	7	5	5909.7285		1	001	11	8	4	6167.6226	0.6	4
100	11	7	4	5909.8120	1.4	2	001	11	8	3	6167.6298		1
100	11	8	4	6100.7209		1	001	11	9	3	6367.8641	0.2	3
100	11	8	3	6100.7245	1.0	4	001	11	9	2	6367.8641	0.5	2
100	11	9	3	6311.2599		1	001	11	11	1	6804.8052		1
100	11	9	2	6311.2594		1	001	11	11	0	6804.8032	0.2	1
100	11	10	2	6545.9418		1	001	12	7	5	6272 5146	0.2	2
100	11	10	1	6545.9421		1	001	12	/ 0	5	6455 6141	0.5	2
100	12	1	10	5566 2226		1	001	12	8	1	6455 6457	0.9	1
100	12	2	10	5560 2001	0.1	1	001	12	9	4	6656 7743	13	2
100	12	3	0	5712 3704	0.1	1	001	12	9	3	6656 7719	0.9	3
100	12	1	0	5726 3086	0.6	1	001	13	5	8	6316 4003	0.4	2
100	12	5	8	5870 8142	0.0	3	001	13	6	8	6418.7810	0.4	2
100	12	5	7	5903.0037	0.0	1	001	13	6	7	6431.2897	0.1	2
100	12	6	6	6030 9186		1	001	13	7	7	6582.9093	0.1	2
100	12	7	6	6196.4569	0.1	2	001	13	7	6	6584.4453	1.1	2
100	12	7	5	6196.7270		1	001	13	8	6	6766.4270	0.8	2
100	12	8	5	6387.6796		1	001	13	8	5	6766.5438		1
100	12	8	4	6387.6944		1	001	14	3	12	6239.6337	0.0	2
100	12	9	4	6599.7013		1	001	14	3	11	6426.9046		1
100	12	9	3	6599.7004		1	001	14	4	11	6429.7520	0.6	3
100	13	1	12	5645.9614	0.7	2	001	14	4	10	6573.8670	0.8	2
100	13	2	12	5644.4435		1	001	14	5	10	6595.3178		1
100	13	2	11	5846.1445		1	001	14	5	9	6683.1843	0.3	2
100	13	3	11	5845.9216		1	001	14	6	9	6751.3228		1
100	13	3	10	6013.4969		1	001	14	6	8	6774.9442		1
100	13	4	10	6021.1418		1	001	14	7	7	6920.3004	0.2	2
100	13	4	9	6135.0728	0.3	2	001	14	8	6	7100.2245		1
100	13	5	9	6176.6638		1	001	15	1	14	6318.4331		1

(continued on next page)

Table 3 (continued)

V         J $K_a$ $K_c$ E $\delta$ 001         15         2         13         6553.1338         0.5							Table 3 (continued)						
V	J	$K_a$	$K_c$	Ε	δ	N	V	J	Ka	$K_c$	Ε	δ	Ν
001	15	2	13	6553.1338	0.5	2	030	6	3	3	5412.8464	0.8	4
001	15	3	13	6554.1362	0.8	2	030	6	4	3	5583.8994	0.4	6
001	15	3	12	6759.1130	1.7	2	030	6	4	2	5584.2591	0.4	2
001	15	4	12	6760.9274	0.7	2	030	6	5	2	5795.1208	0.4	2
001	15	4	11	6928.0323	0.2	1	030	6	5	1	5795.1390	0.3	3
001	15	5	10	6940.9377	0.2	2	030	6	6	1	6034.0149	0.1	2
001	15	5	10	7104 8024	1.5	1	030	07	0	0	5034.0113 5236 1776	0.4	1
001	15	1	10	6633 1669	1.5	1	030	7	1	7	5238.0483	0.4	1
001	16	2	15	6633 1421		1	030	7	1	6	5389 0007	0.2	3
001	16	2	14	6885.5872	0.8	2	030	, 7	2	6	5413.3323	0.2	2
001	16	3	13	7108.2041		1	030	7	2	5	5484.2764		1
001	16	4	13	7109.2586		1	030	7	3	5	5574.5217	0.4	3
001	16	4	12	7297.0275		1	030	7	3	4	5589.9910	0.8	4
001	17	1	17	6663.9213		1	030	7	4	4	5754.4955	0.5	4
001	17	1	16	6965.0865		1	030	7	4	3	5755.7185	0.6	5
001	17	2	16	6965.0608		1	030	7	5	3	5966.0254	1.4	2
001	17	3	15	7234.3224		1	030	7	5	2	5966.1099	0.5	3
001	18	0	18	6995.3560	0.8	2	030	7	6	1	6205.7837		1
001	18	1	18	6995.3568		l	030	7	7	1	6466.6353	0.0	1
001	18	1	17	/314.2251		1	030	/	/	0	6466.6353	0.2	2
001	18	2	1/	/314.22/1		1	030	8	0	8	5391.6628	0.8	1
001	19	1	19	7344.0732		1	030	0	1	0 7	5575 6122	0.8	2
030	19	0	19	4648 4771		1	030	8	2	7	5591 5935	0.2	2
030	1	0	1	4672 2266	0.5	4	030	8	2	6	5691 5536	0.0	2
030	1	1	1	4698.5317	0.2	3	030	8	3	6	5765.6660	0.1	4
030	1	1	0	4704.6302	0.5	3	030	8	3	5	5795.0399	0.1	2
030	2	0	2	4718.7411	0.6	4	030	8	4	5	5949.1204	0.3	4
030	2	1	2	4739.9705	0.4	7	030	8	4	4	5952.7443		1
030	2	1	1	4758.2324	0.1	4	030	8	5	4	6160.9929	0.4	2
030	2	2	1	4834.6627	0.1	5	030	8	5	3	6161.2749		1
030	2	2	0	4835.6053	0.3	5	030	8	6	3	6401.5706		1
030	3	0	3	4786.2349	0.2	4	030	8	6	2	6401.6196		1
030	3	1	3	4801.5276	0.7	5	030	8	7	2	6663.4199		1
030	3	1	2	4837.8491	0.5	6	030	8	8	1	6940.7546		1
030	3	2	2	4906.0978	1.0	8	030	8	8	0	6940.7552	0.0	1
030	3	2	1	4910.6282	0.6	8	030	9	0	9	5564.5200	0.9	2
030	2	2	1	5042.1397	0.4	4	030	9	1	9	5778 0770	0.1	2
030	3 4	0	4	4872 7137	0.7	3	030	9	2	8	5788 9918	0.1	2
030	4	1	4	4882 6337	0.5	7	030	9	2	7	5920 6215	0.5	3
030	4	1	3	4942.4425	0.8	5	030	9	3	7	5978.4992	010	1
030	4	2	3	5000.4339	0.4	7	030	9	3	6	6027.4123	1.2	2
030	4	2	2	5013.0632	0.9	6	030	9	4	5	6175.3114	0.4	3
030	4	3	2	5139.2836	0.5	6	030	9	5	4	6380.5079	1.0	3
030	4	3	1	5139.9384	0.6	5	030	9	6	4	6621.2067		1
030	4	4	1	5315.6843	0.6	5	030	9	6	3	6621.4141		1
030	4	4	0	5315.6940	1.0	3	030	9	7	2	6883.8908		1
030	5	0	5	4976.7851	0.6	5	030	9	8	2	7162.4959		1
030	5	1	5	4982.7216	0.2	3	030	9	8	10	/162.5014		1
030	5	1	4	50/0.49/1	0.9	5	030	10	0	10	5755 1202	0.4	1
030	5	2	4	51/0.9570	0.4	3 7	030	10	1	0	5008 0263	0.4	2
030	5	3	3	5260 6654	1.0	7	030	10	2	9	6005 0197	0.2	2
030	5	3	2	5263 1536	0.6	6	030	10	3	8	6212.0999	0.0	2
030	5	4	2	5437.5894	0.8	4	030	10	4	7	6408.4893		1
030	5	4	1	5437.6654	0.8	6	030	10	4	6	6424.7012		1
030	5	5	1	5648.4248		1	030	10	6	5	6864.5129		1
030	5	5	0	5648.4242	0.6	4	030	11	0	11	5962.6038	0.4	2
030	6	0	6	5097.9364		1	030	11	1	11	5962.7868		1
030	6	1	6	5101.3117	0.8	3	030	11	1	10	6235.6571		1
030	6	1	5	5220.0547	0.2	3	030	11	6	6	7131.4893		1
030	6	2	5	5254.8642	0.9	5	030	11	6	5	7134.5964		1
030	6	2	4	5300.9250	0.8	5	030	12	1	12	6187.9856		1
030	6	3	4	5405.9352	0.6	5	030	12	2	10	6717.1354		1

Table 3	(continued	<i>l</i> )				Table 3 (continued)							
V	J	$K_a$	$K_c$	Ε	δ	N	$\overline{V}$	J	$K_a$	$K_c$	Ε	δ	N
030	12	3	10	6741.2305		1	011	11	7	4	7663.5261	0.3	3
030	12	4	9	6954.6075		1	011	11	8	4	7871.6578		1
030	12	5	8	7183.6783		1	011	11	8	3	7871.6583		1
030	12	6	7	7421.4187		1	011	11	9	3	8096.1766		1
030	13	2	11	7016.9270		1	011	11	9	2	8096.1766		1
030	13	6	8	7733.0660		1	011	12	7	6	7951.6604	0.1	2
030	14	2	12	7332.7516		1	011	12	7	5	7951.9839	0.3	2
030	14	3	12	7339.8673	0.7	2	011	12	8	5	8160.3663		1
030	15	3	12	7895.5905		1	011	12	8	4	8160.3681	0.9	2
110	7	7	1	6684.3091		1	011	12	9	3	8386.1401	0.0	2
110	/		0	6684.3093		1	011	13	3	10	7850 5614	0.8	2
110	8	6	2	6681.6467	0.4	1	011	13	4	9	7830.3014	0.4	2
110	8	7	2	08//.820/	0.4	3	011	13	5	8	/94/.1/28	1.4	2
110	0	/ 0	1	7004 6002	0.8	2	011	13	6	8 7	80/4.//10	1.5	1
110	0	o o	1	7094.0902	0.8	1	011	13	0	7	8262 7641	0.1	2
110	9	6	4	6898 1457	17	2	011	13	7	6	8263 6447	0.1	1
110	0	7	3	7094 9237	1.7	1	011	13	8	6	8471 6543		1
110	9	7	2	7094.9237	0.2	2	011	13	3	12	7845 1412		1
110	9	8	1	7313 2128	0.2	1	011	14	3	11	8043 7642	14	2
110	10	2	8	6654 1277		1	011	14	4	11	8051 1389		-
110	10	3	7	6758.8410	0.8	2	011	14	4	10	8194.0785	0.5	2
110	10	4	6	6844.3094	0.4	3	011	14	5	10	8232.2211		1
110	10	5	5	6970.3475	0.8	3	011	14	5	9	8296.6317		1
110	10	6	5	7138.3040	0.3	2	011	14	6	9	8407.7967		1
110	10	6	4	7138.6208		1	011	14	6	8	8424.7167		1
110	10	7	4	7335.4278		1	011	14	7	7	8598.6700		1
110	10	7	3	7335.4436		1	011	15	1	14	7902.6592		1
110	10	8	3	7555.0455		1	011	15	2	14	7904.6964	1.4	2
110	11	2	10	6727.6103		1	011	15	2	13	8159.6625		1
110	11	3	9	6912.2515		1	011	15	3	13	8161.1956		1
110	11	3	8	7031.9419	0.7	2	011	15	3	12	8380.5136		1
110	11	4	8	7069.0548	0.4	2	011	15	4	12	8384.8069		1
110	11	4	7	7119.7293	0.0	1	011	15	4	11	8554.2435		1
110	11	6	6	7401.9120	1.7	2	011	15	5	11	8579.7038		1
110	11	6	5	7402.8466	0.6	2	011	15	6	10	8762.0454	0.0	1
110	11	/	2	/599.1322		1	011	16	1	16	/900.3385	0.8	2
110	11	/	4	/399.18/4		1	011	16	1	15	8217.2587		1
110	11	ð 1	5 11	/819.9198		1	011	10	2	15	8402 4062		1
110	12	1	11	6073 6685	0.0	2	011	16	2	14	8493.4902		1
110	12	2	10	7170 2513	0.0	1	011	16	3	14	8732 9250		1
110	12	4	9	7346 8085		1	011	17	0	17	8211 4060		1
110	12	5	7	7534 7083		1	011	17	1	17	8211 4070		1
110	12	6	6	7691.1132		1	011	17	2	16	8548.8697		1
110	12	7	6	7885.8216		1	011	18	0	18	8539.6131		1
110	13	0	13	6974.4174	1.1	2							
110	13	1	13	6974.5341		1							
110	13	1	12	7233.7194		1							
110	13	5	9	7815.6391		1							
110	14	0	14	7232.7676		1	not g	given in	our l	inelist (	Supplemental	Materi	ial I)
110	14	1	14	7232.7672		1	becaus	se they a	are too s	strong of	· blended by a	diacent	lines.
110	14	2	13	7510.6075		1	On th	e other	hand	totally '	7341 new trai	reitione	were
110	15	0	15	7508.1380		1	ohoom	rad In	manu,	1000000000000000000000000000000000000	consistions of a	ha 2.	hand
110	15	1	15	7508.1374		1	observ			11 330 L		$110 SV_2$	band
011	9	9	1	7584.6735	0.9	2	were	obtained	I. Morec	over, 93	lines show la	rge disa	agree-
011	9	9	0	7584.6723	0.0	2	ment	between	the va	lues in	HITRAN2004	and in	n this
011	10	5	6	7045.0169	0.1	2	work,	87 of th	nem are	listed in	HITRAN2004	but no	ot evi-
011	10	7	4	7398.3291	0.1	2	denced	d in ou	r spectru	ım. The	largest deviat	ion is	up to
011	10	7	3	/398.3547	0.8	5	-930	$m^{-1}$ . It	implies	that so	me mis-assion	ed or r	voorly
011	10	8	3	/003.8516	0./	4	simula	ted li	nec in	this	region remo	in in	the
011	10	ð	2	7003.8319	0.0	4			ldotob-	00 The-			,
011	10	9	2 1	7828 0127	1.0	2	ппк	AIN2004	uataba	se. Thes	e transitions	ogether	with
011	10	9 10	1	8063 7008	0.0	ے 1	our as	signmen	its can b	e tound	in the Supple	mental	Mate-
011	10	10	0	8063 7998		1	rial I	attached	to this	paper.			
	10	10	0	00000.1990		1							

In the PS calculation, 8824 transitions of  $H_2^{18}O$  are predicted to be stronger than  $1.0 \times 10^{-6} \text{ cm}^{-1}/(\text{atm cm})$ in the studied region including hot transitions. The data of HITRAN2004 database in the studied region comes from Refs. [12,17–19,22]. Owing to relatively longer

Table 4				
H <sub>2</sub> <sup>18</sup> O energy levels	which is much	different with	those in	[34]

V	J	Ka	$K_c$	This work			Ref. [34]	
				Ε	δ	N	Ε	δ
010	10	9	2	4207.2975	0.64	4	4208.520220	0.4
010	10	9	1	4207.2970		1	4208.520220	0.4
010	10	10	1	4458.5028	1.03	2	4467.825950	3.0
010	10	10	0	4458.5016	0.07	2	4467.825950	3.0
010	11	6	5	3812.6606	0.23	4	3812.680000	3.0
010	11	9	3	4476.9332		1	4477.870600	3.0
010	11	9	2	4476.9331	0.59	3	4477.870600	3.0
010	12	4	8	3831.1147	0.68	5	3831.070000	3.0
010	12	6	7	4102.4117	0.28	4	4102.043100	4.0
010	12	6	6	4105.0445	0.58	5	4104.731430	4.0
010	12	7	6	4304.3001	0.42	5	4303.652600	2.0
010	12	7	5	4304.5020	0.63	3	4303.866220	4.0
010	12	8	5	4528.0136	0.95	2	4527.266900	4.0
010	13	4	9	4161.1780	0.02	3	4160.835570	2.0
010	13	5	9	4234.4176	0.34	3	4234.005880	4.0
010	13	5	8	4270.2918	0.59	5	4270.018360	2.0
010	13	6	7	4422.2331	0.44	4	4421.703200	3.0
010	14	2	12	4166.6145	0.20	2	4166.661910	3.0
010	14	3	11	4366.8762	0.49	2	4366.483400	3.0
010	14	5	10	4566.5879	0.54	2	4565.874270	4.0
010	15	3	13	4489.1903		1	4489.390050	4.0
010	15	3	12	4710.5253	0.91	2	4710.032710	4.0
010	16	1	15	4545.1661		1	4545.295700	4.0
010	16	2	15	4545.1537		2	4545.313200	4.0
010	18	0	18	4869.1798		1	4869.249500	4.0
010	18	1	18	4869.1798		1	4869.249500	4.0
020	7	5	3	4345.1004	0.53	5	4345.138720	0.4
020	13	1	13	4921.6103		1	4921.599000	3.0
100	10	7	4	5646.2033	0.42	3	5646.200880	3.0
100	11	2	9	5303.2254		1	5303.231280	0.8
100	11	3	8	5426.8684	0.61	3	5426.864710	1.0
001	8	7	2	5261.0099		1	5261.003740	1.0
001	12	1	12	5268.4148		1	5268.430300	3.0
001	13	4	9	6235.1536	0.22	3	6235.079880	0.4
110	3	2	2	5433.5423	0.58	5	5433.550180	0.2
110	4	4	0	5736.5663	0.40	2	5736.577300	1.0
110	6	5	2	6147.7326	0.36	6	6147.735200	1.5
110	8	5	4	6509.1524	0.71	6	6509.147450	1.5
110	10	5	6	6966.6160	0.71	5	6966.611540	1.5
110	11	5	7	7232.1400	0.72	3	7232.137800	2.0
110	11	5	6	7239.2275		1	7239.378400	5.0
011	9	7	2	7156.5940	0.80	2	7156.588300	0.4
011	12	1	12	6829.2725		1	6829.277900	3.0
011	12	4	8	7527.2949	0.18	3	7527.291470	1.5
011	12	5	8	7596.1095	0.59	3	7596.104500	2.0
011	12	6	6	7767.5980	0.43	3	7767.611000	2.0
011	13	2	11	7545.1740	0.68	2	7545.178120	1.5
011	14	1	14	7328.0896	0.73	2	7330.408100	3.0
011	14	1	13	7606.5469	1.00	4	7606.542530	1.5
011	14	2	13	7606.1635	0.47	2	7606.157500	2.0
011	16	0	16	7900.3387	1.02	2	7900.286400	2.0

*Note.* N is the number of lines used for the upper level determination and  $\delta$  denotes the experimental uncertainties (combination differences) in  $10^{-3}$  cm<sup>-1</sup>.

absorption path length and higher partial pressure of the  $H_2^{18}O$  sample in our experiment, we observed more weak lines. While in some regions, many strong lines are saturated. Together with previous work, totally 6534 transitions have been observed. With very careful comparison with the observations and predictions, we think the remainder are blended by other strong transitions of  $H_2^{18}O$  or other isotopologues of water. Fig. 3 shows the stick spectrum in 4000–6000 cm<sup>-1</sup> region provided by this work, HITRAN2004, and PS calculation. Obviously there are much more transitions contained in this work than in HITRAN2004, and our results agree well with the PS calculations.

Table 2 gives a summary of H<sub>2</sub><sup>18</sup>O energy levels determined in this work. The result given by [34] is also presented as a comparison. In this work, totally 1339 rotational energy levels of 9 different upper vibrational states were determined, while the number is 432 from [34]. In particular, the rotational energy levels of the (030) state are mostly obtained for the first time. The complete list of the energy levels is given in Supplementary Material II attached to this paper. Only the newly derived energy levels are given in Table 3 here. Totally 946 energy levels were confirmed through combination differences of two or more transitions. The average experimental uncertainty for these levels was estimated to be about  $0.0005 \text{ cm}^{-1}$ . According to the systematic tendency of the Eobs-calc values (difference between observed and PS predicted energies) and the line intensity predictions, energy levels derived from one transition can be validated, but may have considerably larger uncertainty. Besides of 432 new energy levels, 51 energy levels are very different from those in [34]. For these very different levels, 44 of them given in [34] have large uncertainties (up to  $0.03 \text{ cm}^{-1}$ ), some rise up from transitions do not appear in our spectrum, which could be in error. For our derivation of such levels, they are confirmed by the comparison with the PS calculations (both line positions and intensities), and 39 of them even with good combination differ-

Table 5 Summary of the determined energy levels of HD<sup>18</sup>O

V	$E_v (\mathrm{cm}^{-1})$	$J^{\max}$	$K_a^{\max}$	Number of	Levels		Number of
_				This work	Ref. [34]	New	Transitions
010	1396.267	18	9	182	137	47	807
100	2709.284	14	7	127	105	24	632
020	2767.212	12	6	91	59	32	500
001	3696.330	17	9	163	128	36	891
110	4080.545	14	6	101	83	18	458
030	4121.754	12	4	76	55	21	331
011	5071.497	14	8	121	0	121	891
200	5335.361	13	5	77	0	77	172
002	7229.186	14	7	128	0	128	523
Total				1066	567	504	4942

Table 6

Table 6	(continued)

New H (001) st	tew HD <sup>18</sup> O ro-vibrational energy levels of the (010), (020), (100), 001) states (cm <sup><math>-1</math></sup> )						h, and $\frac{1 \text{ able 6 (continuea)}}{V  J  K_a  K_c  E  \delta}$							
<u>V</u>	I	) K	K	F	δ	N	020	6	6	1	3765.5563		1	
<i>v</i>	J	$\mathbf{\Lambda}_{a}$	$\mathbf{\Lambda}_{\mathcal{C}}$	L	0	11	020	7	4	3	3513.7604	0.8	4	
010	8	8	0	2991.1921	0.1	2	020	7	4	4	3513.5185	0.4	6	
010	8	8	1	2991.1921	0.1	2	020	7	5	2	3678.6335		1	
010	9	8	1	3128.7208	0.4	2	020	7	5	3	3678.6323		1	
010	9	8	2	3128.7208	0.4	2	020	8	2	6	3440.6608*	0.4	4	
010	9	9	0	3386.0115		1	020	8	3	5	3515.8779	0.6	5	
010	9	9	1	3386.0115		1	020	8	3	6	3506.1184	0.6	5	
010	10	7	3	3050.7079	1.7	2	020	8	4	4	3639.5768	0.3	4	
010	10	/	4	3050.7093	0.1	2	020	8	4	5	3638.8755	0.5	3	
010	10	8	2	3281.5131	0.1	3	020	9	2	7	3591.0107	0.8	5	
010	10	8	3	3281.5135	0.4	3	020	9	2	8	3526.7059*	0.6	6	
010	10	9	1	3537.9484		1	020	9	3	6	3663.7908	0.7	5	
010	10	9	2	3337.9484	1 1	1	020	9	3	7	3646.2080	0.7	5	
010	11	0	5	2017.0810	1.1	2	020	9	4	5	3781.7934	1.0	5	
010	11	07	0	2210 7276	0.0	2	020	9	4	6	3780.0460	0.7	3	
010	11	7	4	3219.7370		1	020	10	1	9	3665.3693	0.5	3	
010	11	/ Q	2	3219.7374		1	020	10	2	8	3756.3812	0.2	2	
010	11	0	5	2449.3420		1	020	10	2	9	3672.9135	0.6	4	
010	11	0	4	2816 5710	1.2	1	020	10	3	7	3829.6595		1	
010	12	3	9	2010.3719	0.7	2	020	10	3	8	3801.0517	1.0	4	
010	12	4	0	2901.2141	0.7	4	020	11	0	11	3672.5174	0.4	4	
010	12	4	9 7	20032 0004	0.0	4	020	11	1	10	3827.3595	0.3	2	
010	12	5	/ 0	2020 2640	1.5	2	020	11	1	11	3672.7066	0.9	4	
010	12	5	6	3202 8245	0.7	2	020	11	2	10	3832.0507	1.0	4	
010	12	7	5	3404 1541	0.9	1	020	12	0	12	3829.1443	1.1	2	
010	12	7	6	3404 1531		1	020	12	1	12	3829.2428	0.3	2	
010	12	3	10	3030 7361	0.9	4	100	6	5	2	3402.4632*	0.8	3	
010	13	3	11	2951 8042	0.7	3	100	8	7	1	3976.9806	0.1	2	
010	13	4	9	3113 0114	0.9	3	100	8	1	2	3976.9806	0.1	2	
010	13	4	10	3087 6435	0.5	2	100	9	6	3	3924.7221	1.2	3	
010	13	5	9	3233 0517	0.1	2	100	9	6	4	3924.7211	1.8	4	
010	14	1	13	2999 1195*	0.5	4	100	10	I	9	35/3.8685*	0.9	6	
010	14	2	12	3152 3001	0.1	1	100	10	4	/	3/89./069*	0.2	3	
010	14	2	13	2999 7157*	1.0	3	100	10	5	2	3919.2559	0.2	3	
010	14	3	11	3259.1276		1	100	10	5	0	3918.9/10	1.0	2	
010	14	3	12	3160.5271	0.3	2	100	10	6	4	40/5.840/	1.1	2	
010	14	4	10	3342.4247		1	100	11	2	9	3624.4797	1.5	4	
010	14	4	11	3304.0607		1	100	11	2	10	3734.0764	0.0	4	
010	15	1	14	3500.6075		1	100	11	2	0	3887.9070	0.8	4	
010	15	2	14	3208.3146	0.5	2	100	11	3	7	3040.7002	1.1	2	
010	15	3	12	3207.9880	0.5	2	100	11	4	8	3957 7405	0.8	3	
010	15	3	13	3382.0060		1	100	11		6	4087 2614	0.9	1	
010	15	4	11	3588.7506		1	100	11	5	7	4086 5306	0.0	2	
010	16	1	15	3754.1234		1	100	12	1	11	3900 9751	0.0	4	
010	17	0	17	3431.5413	1.2	2	100	12	2	10	4009 5552	0.0	4	
010	17	1	17	3431.5413	1.2	2	100	12	2	11	3902 4388	0.3	3	
010	17	2	15	4128.1632		1	100	12	3	9	4081 5967	0.3	3	
010	18	0	18	3662.4399	0.6	2	100	12	3	10	4025 0022	0.5	3	
010	18	1	18	3662.4399	0.6	2	100	13	0	13	3927 6035*	11	3	
020	4	2	3	3000.9061*	0.6	7	100	13	1	13	3927.6035*	1.1	3	
020	5	2	4	3077.5223*	0.5	10	100	13	2	11	4206.3337	0.4	2	
020	5	4	1	3310.2743*	0.5	4	100	13	2	12	4083.2118	0.9	3	
020	5	4	2	3310.2606*	1.1	4	100	14	0	14	4107.9437		1	
020	5	5	0	3476.5602	1.5	2	100	14	1	14	4107.9608		1	
020	5	5	1	3476.5574	0.1	2	001	10	3	8	4671.2933*	1.0	6	
020	6	2	4	3188.6710*	0.4	6	001	10	6	4	5044.4824	0.3	2	
020	6	3	4	3271.5497*	0.6	7	001	10	6	5	5044.4647	1.1	3	
020	6	4	2	3404.0529	1.2	3	001	10	7	3	5217.9254	0.4	2	
020	6	4	3	3403.9936	0.4	5	001	10	7	4	5217.9276	1.8	3	
020	6	5	1	3569.7720	0.7	2	001	10	8	2	5415.9356	1.6	2	
020	6	5	2	3569.7691	0.6	2	001	10	8	3	5415.9357	1.6	2	
020	6	6	0	3765.5563		1	001	10	9	1	5637.4394	1.4	2	

(continued on next page)

Table 6 (continued)

V	J	Ka	$K_c$	Ε	δ	N
001	10	9	2	5637.4394	1.4	2
001	11	4	7	4956.9408	1.5	4
001	11	4	8	4945.2065		1
001	11	5	6	5068.2171	0.7	3
001	11	5	7	5067.0483		1
001	11	6	5	5213.6119		1
001	11	6	6	5213.5497	0.0	2
001	11	7	4	5385.9305		1
001	11	7	5	5385.9303		1
001	12	1	11	4893.3999*	0.0	2
001	12	2	10	5006.2227	1.0	4
001	12	3	9	5081.7725		1
001	12	3	10	5017.2449	1.6	3
001	12	4	8	5150.7580	1.0	3
001	12	4	9	5130.5093	1.7	2
001	12	5	7	5255.5903		1
001	13	2	11	5203.2065	1.1	2
001	13	3	10	5075.4038	0.7	2
001	13	3	11	5210.0618	1.5	3
001	13	4	9	5362.1813		1
001	13	4	10	5330.3069		1
001	14	2	12	5411.5676	1.2	2
001	14	2	13	5269.7835		1
001	14	3	11	5269.5363		1
001	14	3	12	5415.6432	0.7	2
001	15	0	15	5294.8465*	0.5	2
001	15	1	14	5752.1256		1
001	15	1	15	5294.8477*	0.8	2
001	16	0	16	5500.5313	0.6	3
001	16	1	16	5500.5313	0.6	3
001	17	0	17	5718.3829	0.3	2
001	17	1	17	5718.3829	0.3	2

*Note.* N is the number of lines used for the upper level determination and  $\delta$  denotes the experimental uncertainties (combination differences) in  $10^{-3}$  cm<sup>-1</sup>. Energy levels different from [34] are marked by '\*'.

ences. Table 4 lists such energy levels which are very different with those in [34].

#### 3.2. HD<sup>18</sup>O

The isotopic abundance of HD<sup>18</sup>O molecule in the measurement with 1307 Pa sample was estimated to be about 0.8% by comparing the experimental line intensities of some moderate lines with the value given in [29]. In the 1080–7800 cm<sup>-1</sup> region, there are 9 bands of HD<sup>18</sup>O:  $v_2$ ,  $v_1$ ,  $2v_2$ ,  $v_3$ ,  $v_1 + v_2$ ,  $3v_2$ ,  $v_2 + v_3$ ,  $2v_1$ , and  $2v_3$ . There are also many lines due to H<sub>2</sub><sup>18</sup>O, H<sub>2</sub><sup>16</sup>O, and HD<sup>16</sup>O in this region. It makes the assignments of HD<sup>18</sup>O complicated. The line positions of H<sub>2</sub><sup>16</sup>O and HD<sup>16</sup>O can be known from HITRAN database and our previous measurements of the water sample. The new results of H<sub>2</sub><sup>18</sup>O presented in this paper are also used to identify the lines. All bands of HD<sup>18</sup>O studied in this work are found to be of the A/B hybrid type.

First, the experimental ground-state rotational energy levels of [34] were fitted with the Watson type Hamiltonian [35]:

$$\begin{aligned} H^{v} &= E^{v} + \left[A^{v} - \frac{1}{2}(B^{v} + C^{v})\right]J_{z}^{2} + \frac{1}{2}(B^{v} + C^{v})J^{2} + \frac{1}{2}(B^{v} - C^{v})J_{xy}^{2} \\ &- \Delta_{K}^{v}J_{z}^{4} - \Delta_{JK}^{v}J_{z}^{2}J^{2} - \Delta_{J}^{v}J^{4} - \delta_{K}^{v}\left[J_{z}^{2}, J_{xy}^{2}\right]_{+} - 2\delta_{J}^{v}J^{2}J_{xy}^{2} + H_{K}^{v}J_{z}^{6} \\ &+ H_{KJ}^{v}J_{z}^{4}J^{2} + H_{JK}^{v}J_{z}^{2}J^{4} + H_{JJ}^{v}J^{6} + \left[J_{xy}^{2}, h_{K}^{v}J_{z}^{4} + h_{JK}^{v}J^{2}J_{z}^{2} + h_{J}^{v}J^{4}\right]_{+} \\ &+ L_{K}^{v}J_{z}^{8} + L_{KKJ}^{v}J_{z}^{6}J^{2} + L_{JK}^{v}J_{z}^{4}J^{4} + L_{KJJ}^{v}J_{z}^{2}J^{6} + L_{J}^{v}J^{8} \\ &+ \left[J_{xy}^{2}, l_{K}^{v}J_{z}^{6} + l_{KJ}^{v}J^{2}J_{z}^{4} + l_{JK}^{v}J^{4}J_{z}^{2} + l_{J}^{v}J^{6}\right]_{+} \\ &+ P_{K}^{v}J_{z}^{10} + \left[J_{xy}^{2}, p_{K}^{v}J_{z}^{8}\right]_{+} + \dots, \end{aligned}$$

where  $J_{xy}^2 = J_x^2 - J_y^2$  and  $[A, B]_+ = AB + BA$ . The ground state constants were obtained with  $1.95 \times 10^{-4}$  cm<sup>-1</sup> standard deviation, which can reproduce the experimental values pretty well. So a set of the ground state rotational energy levels was formed by calculated values. Actually, the analysis of  $v_2$ ,  $v_1$ ,  $v_3$ ,  $2v_2$ ,  $3v_2$ , and  $v_1 + v_2$  bands were almost straightforward based on the previous work of Toth [34]. More assignments and energy levels have been obtained and used to refine the spectroscopic parameters. The  $v_2 + v_3$ ,  $2v_1$ , and  $2v_3$  bands are observed for the first time, and they are strongly mixed with the absorptions of other isotopes of water. So these three bands were analyzed simultaneously with the determination of the Hamiltonian parameters in the "snow ball" method [36]. Table 5 gives a summary of assignments and energy levels obtained in this work and [34].

From 4942 transitions, 1062 final rotational energy levels of nine upper states were determined. The whole list can be found in Supplemental Material II. Tables 6 and 7 only give new derived energy levels. Rotational energy levels of  $3v_2$  and  $v_1 + v_2$  states are all listed for convenience. Most energy levels were determined by two or more transitions from different ground state levels.  $\delta$  (in  $10^{-3}$  cm<sup>-1</sup>), the experimental uncertainties which come from combination differences, and N, the number of lines used for the upper level determination are also presented in the same table. For most levels,  $\delta$  values are less than 0.0012 cm<sup>-1</sup>, which represent the experimental accuracy. No transitions to 28 rotational energy levels in [34] are evidenced in our spectrum, 25 of them are replaced in this work. The ro-vibrational parameters of seven vibrational states (except (030) and (110)) were fitted by nonlinear least-squares method with Watson type Hamiltonian. The parameters are listed in Table 8. It is clear that these seven vibrational states can be treated well as isolated ones.

#### 4. Conclusion

The high resolution Fourier-transform spectrum of <sup>18</sup>O enriched water vapor was recorded and analyzed in the 1080–7800 cm<sup>-1</sup> spectral region. Analysis of the  $3v_2$  band of H<sub>2</sub><sup>18</sup>O and the  $v_2 + v_3$ ,  $2v_1$ , and  $2v_3$  bands of HD<sup>18</sup>O are reported for the first time. Totally 5231 and 4942 transitions were assigned to H<sub>2</sub><sup>18</sup>O and HD<sup>18</sup>O, respectively. A set of 936 new accurate energy levels was derived for these two isopologues of water molecule.

Table 7 Full list of the ro-vibrational energy levels of the (110), (030), (011), (200), and (002) states of  $HD^{18}O$  (in cm<sup>-1</sup>)

		(002)			(200)			(011)			(030)			(110)			
J	$K_a$	$K_c$	Ε	δ	N	Ε	δ	N	Ε	δ	N	Ε	δ	N	Ε	δ	N
	1		2	3	4	5	6	7	8	9	10	11	12	13	14	15	16
0	0	0	7229.1857		1	5335.3609		1	5071.4966	0.5	2	4121.7545	0.3	2	4080.5450		1
1	0	1	7244.4294	0.8	4	5350.2167	0.1	2	5086.8842	0.2	3	4137.1960	0.1	3	4095.8601	0.5	4
1	1	0	7259.2981	0.9	4	5366.6667		1	5104.6327	0.3	4	4159.5250	0.2	4	4117.3437	0.2	2
1	1	1	7256.4721	0.3	3	5364.1265		1	5101.6566	0.2	5	4156.4366	0.3	3	4114.3068	0.4	4
2	0	2	7274.4659	0.6	5	5379.6046	0.3	2	5117.2475	0.3	5	4167.7338	0.3	3	4126.1276	0.2	4
2	1	1	7292.5943	0.4	4	5398.9048	0.3	3	5138.3678	0.2	5	4193.5463	0.2	5	4150.9246	0.2	4
2	1	2	7284.1234	0.3	6	5391.2817		1	5129.4476	0.3	6	4184.2795	0.8	5	4141.8354	0.7	5
2	2	0	7329.0392	0.6	4	5440.7292		1	5182.8407	0.6	6	4253.5377	0.4	6	4203.4486	0.7	5
2	2	1	7328.5998	0.2	5	5440.4140		1	5182.4367	0.3	4	4253.2000	0.9	7	4203.0983	0.2	7
3	0	3	7318.4706	0.8	4	5422.9049	0.7	3	5161.8093	0.8	8	4212.7041	0.6	6	4170.6566	0.6	3
3	1	2	7342.2116	0.9	6	5447.0262	0.1	3	5188.6697	0.2	7	4244.3489	0.4	6	4201.0068	0.4	5
3	1	3	7325.3315	0.6	7	5431.8105	0.1	3	5170.8847	0.7	9	4225.8312	0.7	7	4182.9113	0.6	5
3	2	1	7376.3983	0.2	4	5486.4534	0.4	3	5230.5062	0.3	7	4301.3859	0.7	6	4250.4318	0.5	8
3	2	2	7374.2706	0.6	8	5484.9054		1	5228.5397	0.3	8	4299.7299	0.8	7	4248.7154	0.8	8
3	3	0	7441.2291	0.4	4	5560.0197	0.1	2	5308.8963	0.2	5	4409.3005	0.3	3	4337.5020	0.4	4
3	3	1	7441.1858	0.4	4	5559.9948	0.2	2	5308.8619	0.2	4	4409.2786	0.5	4	4337.4753	0.7	1
4	0	4	7375.4988	0.8	6	5479.3528	0.2	3	5219.6322	0.5	7	42/1.2605	0.5	6	4228.5611	0.8	6
4	1	3	/40/.6996	0.9	5	5510./162	0.7	3	5255.1316	0.6	9	4311.6240	0.3	6	4267.2129	0.5	2
4	1	4	7379.8350	0.9	5	5485.5200	1.1	3	5225.7184	0.4	9	4280.8693	0.6	3	4237.3111	0.6	6
4	2	2	7440.7693	1.0	5	5548.3836	1.0	2	5295.2587	0.5	8	4366.2700	0.6	2	4314.15/6	0.7	9
4	2	3	7434.7912	0.2	5	5(10.0400	0.7	2	5289.6641	0.6	6	4361.4804	0.5	1	4309.2303	0.8	/
4	3	1	/503.050/	1.0	6	5619.9408	0.7	3	53/1.1831	0.7	6	44/2.11/3	0.2	4	4398.8585	0.6	9
4	3	2	/502./533	0.6	0	5619.7692	0.4	2	5370.9452	0.5	2	44/1.9642	0.8	4	4398.6734	0.9	0
4	4	0	/594.5943	1.0	3	5723.2047		1	5480.9008	0.6	3	4622.7618	1.2	3	4516.9900	0.7	3
4	4	1	7394.3948	0.5	4	5725.2007	0.1	1	5480.8997	0.3	3	4022.7013	1.0	4	4310.9878	1.2	2
5	0	2	/444.8/8/	0.3	4	5548.2706	0.1	2	5289.9540	0.7	/	4342.614/	0.9	0	4299.0304	0.7	3
5	1	4	7488.3707	0.5	4	5552 1900	0.0	2	5357.1405	0.0	8	4394.9038	0.1	4	4348.9913	0.4	5
5	1	2	7522 4570	0.9	3	5552.1890	0.0	3	5295.0824	0.8	10	4349.1400	0.3	5	4304./890	0.4	5
5	2	3	7522.4570	0.5	4	5617 2145	1.1	4	5265 5222	0.4	3 7	4448./313	0.2	5	4393.0833	0.5	07
5	2	4	7509.8514	0.5	5	5605 2570	0.7	2	5440 5976	0.0	0	4436.2062	1.0	5	4364.369/	0.7	0
5	3	2	7570 8115	0.0	3 7	5604 5846	0.0	5	5449.5870	0.5	9	4550 4222	0.2	5	44/0.005/	0.4	0
5	3	3	7671 4405	0.9	5	5707 8100		1	5558 2507	0.7	7	4330.4323	0.2	1	4475.2625	0.3	0 5
5	4	2	7671 4191	0.0	7	5/9/.0109		1	5558 3389	0.8	6	4700.9903	0.0	4	4593.2523	0.4	6
5	-+ -5	0	7788 1610	0.5	2	5930 0345	18	2	5607 3013	0.7	3	4700.9795	0.4	4	4393.2377	0.0	3
5	5	1	7788 1580	0.8	3	5930.0345	1.0	1	5607 3000	0.0	1				4741.2479	0.8	1
6	0	6	7526 3535	0.5	2	5629 2814	0.9	2	5372 4003	0.4	10	4426 2444	0.6	5	4381 5609	0.7	3
6	1	5	7520.3335	0.9	6	5682 7213	0.5	3	5433 8549	0.4	9	4493 5496	1.0	5	4445 5929	0.4	4
6	1	6	7527 6770	0.9	7	5631 6038	0.0	3	5374 5303	0.5	8	4430 4204	0.2	4	4385 1105	0.0	5
6	2	4	7621 2403	0.8	6	5722 1584	0.9	2	5477 2171	0.6	8	4548 9668	0.5	6	4493 2736	0.5	5
6	2	5	7599.0856	0.5	5	5704.7525	0.6	4	5455.7725	0.7	11	4529.6233	0.4	7	4473.8875	0.7	7
6	3	3	7675.5403	0.7	8	5786.3555		1	5544.6087	0.4	6	4646.4204	0.8	4	4569.3534	0.3	7
6	3	4	7672.2697	0.2	5	5784.4114	1.4	2	5541.9435	0.4	8	4644.6622	0.6	5	4567.2640	0.4	8
6	4	2	7763.9643	0.6	4				5651.5549	0.9	7	4795.0290	0.5	5	4684.9712	0.3	3
6	4	3	7763.8127	0.6	7				5651.4477	0.5	7	4794.9750	0.4	5	4684,8970	0.4	3
6	5	1	7880.0637	0.8	2				5789.9176	0.9	4				4832.3919	0.4	3
6	5	2	7880.0643	0.1	3				5789.9179	0.5	3				4832.3924	0.2	4
6	6	0	8020.9517		1				5956.5304	0.4	3				5009.2343	0.1	3
6	6	1	8020.9494	1.0	4				5956.5299	1.0	3				5009.2343	0.2	3
7	0	7	7619.9225	0.6	5	5722.2772	0.2	2	5466.9039	0.5	7	4521.9294	0.5	5	4475.9838	0.9	3
7	1	6	7691.4392	0.1	2	5789.5644	1.2	3	5544.2587	0.6	8	4606.7313	0.8	3	4556.0881	0.7	6
7	1	7	7620.5940	1.0	3	5723.5830	0.6	3	5468.0606	0.2	4	4524.4786	0.2	4	4478.0671	0.5	5
7	2	5	7736.4991	0.4	4	5833.6951	0.2	2	5593.8679	0.7	7	4666.7652	0.2	3	4608.3747	0.3	6
7	2	6	7702.1022	0.4	6	5805.9411	0.4	2	5560.0391	0.2	7	4635.3939	0.6	3	4577.3834	0.5	6
7	3	4	7787.4046	0.5	5	5893.7360	0.0	2	5656.8531	0.8	7	4758.7571	0.1	4	4679.4380	0.8	5
7	3	5	7779.9240	0.3	8	5889.1302	0.4	3	5650.6370	0.8	6	4754.5552	0.6	5	4674.5002	0.6	6
7	4	3	7872.3677	0.9	7				5760.6655	0.3	3	4904.9863	0.6	4	4792.2887	0.4	6
7	4	4	7871.8266	0.7	7				5760.2855	0.2	7	4904.7964	0.9	5	4792.0231	0.6	4
7	5	2	7987.4607		1				5898.1179	1.0	4				4938.8468	1.1	3

*Note.* N is the number of lines used for the upper level determination and  $\delta$  denotes the experimental uncertainties (combination differences) in  $10^{-3}$  cm<sup>-1</sup>. (*continued on next page*)

Table 7 (continued)

(002)					(200)			(011)			(030)		(110)				
J	<i>K</i> <sub>a</sub> 1	K <sub>c</sub>	E 2	δ 3	N 4	E 5	$\delta$ 6	N 7	E 8	δ 9	N 10	Е 11	δ 12	N 13	Е 14	δ 15	N 16
7	5	3	7987.4466	0.7	2				5898.1060	0.5	2				4938.8435	1.2	5
7	6	1	8127.7872	1.0	3				6064.1724	0.3	4				5114.4963	0.9	2
7	6	2	8127.7875	1.0	3				6064.1720	0.5	4				5114.4973		1
7	7	0	8291.9131	0.9	2				6257.0593		1						
7	7	1	8291.9131	0.9	2				6257.0593		1						
8	0	8	7725.6519	0.5	5	5827.2789	0.1	2	5573.5157	0.5	4	4629.6390	0.5	2	4582.3191		1
8	1	7	7811.8601	1.0	4	5909.1773	0.7	3	5667.3594	0.5	7	4733.5278	0.8	4	4679.4913	0.6	5
8	1	8	7725.9819	0.8	4	5827.9865	0.6	2	5574.1238	0.7	4	4631.1448	0.5	2	4583.5001	0.2	3
8	2	6	7867.3641	0.6	6	5960.9826	0.2	2	5726.7089	0.8	5	4801.6134	1.2	4	4739.7526	0.7	5
8	2	7	7818.5091	0.6	4	5920.5561	0.9	3	5677.9342	0.4	5	4755.1684	0.1	3	4694.5187	0.6	4
8	3	5	/916.94/0	0.9	6	601/.8519		1	5/86.8338	0.4	2	4888.5817	0.8	4	4806.7791	0.6	5
ð	3	0	7902.4632	0.9	0	6008.3327		1	5005 0614	0.0	8	48/9.9308	0.5	2	4/90.//48	0.4	0 5
0	4	4	7997.0004	1.0	4				5884 8666	0.9	4	5030.4682	0.0	3	4913.4100	0.4	5
0 8	+ 5	3	8110 4411	1.0	4				6021 0680	0.8	1	5050.4082	1.0	5	5060 6731	0.0	4
8	5	4	8110.3722	0.6	- -				6021.9089	0.8	3				5060.6471	0.9	4
8	6	2	8250 0046	1.0	3				6187 2765	1.0	2				5000.0471	0.4	7
8	6	3	8250.0024	0.4	3				6187 2751	0.4	2						
8	7	1	8413.5291	0.3	3				0107.2751	0.1	-						
8	7	2	8413.5291	0.4	3												
8	8	0							6598.0065		1						
8	8	1							6598.0065		1						
9	0	9	7843.5895	0.7	3	5944.3315	1.4	2	5692.2957	0.3	4	4749.4123	1.0	3	4700.6351	1.4	2
9	1	8	7944.0807	0.6	5	6040.8649	0.9	3	5802.4531	0.6	3	4873.0725	0.6	2	4814.9792	0.7	5
9	1	9	7843.7469	0.1	2	5944.7077	0.7	3	5692.6074	0.2	4	4750.2821	0.6	3	4701.2868	1.1	2
9	2	7	8012.7785	1.0	5	6103.2614		1	5874.7945	0.6	4	4952.7690	0.4	3	4886.5843	0.1	2
9	2	8	7947.9517	0.9	4	6048.2800	0.1	2	5809.0919	0.8	5	4888.5951	0.8	5	4824.9438	0.7	4
9	3	6	8064.0610	0.4	5				5934.7031	0.5	4	5036.3236	1.0	3	4951.6630	0.5	5
9	3	7	8039.4970	0.1	4	6142.4117	1.3	2	5913.1203	0.8	8	5020.5274	0.4	2	4933.8181	0.2	4
9	4	5	8138.3699	1.0	5				6027.8433	0.8	2				5054.6645	1.0	4
9	4	6	8134.6165	0.6	2				6025.1429	1.0	4				5052.7421	0.4	3
9	5	4	8249.1350	1.0	3				6161.5615	1.0	1				5197.9393	0.7	3
9	5	5	8248.9021	1.0	2				6161.4210	1.0	3				5197.8495	1.0	4
9	6	3	838/.0351	2.0	2				6325.8772	1.0	2						
9	0	4	8550 4121	0.7	3				0323.8770	2.0	2						
9	7	2	8550.4131	0.6	2												
10	0	10	7973 7559	0.0	1	6073 4724	1.0	2	5823 2816	03	4	4881 2947	0.4	2	4830 9901	1.0	2
10	1	9	8087 9937	1.0	3	0073.4724	1.0	4	5949 2381	0.5	4	5024 7181	0.4	2	4962 0782	0.2	3
10	1	10	7973 8246	1.0	2	6073 6663	1.1	2	5823 4390	0.7	4	4881 7875	0.1	$\frac{2}{2}$	4831 3464	0.1	2
10	2	8	8171 5725	0.7	4	6259 6662		1	6037 0318	0.7	4	10011/0/5	0.5	-	5047 9061	0.5	2
10	2	9	8090.1325	0.9	3	6188.8287	0.1	2	5953.1947	1.0	3	5035.3390	0.8	2	4968.3399	0.4	2
10	3	7	8228.1329	0.7	4				6100.1140	0.9	3				5113.9830	0.9	3
10	3	8	8190.5855	0.8	4	6290.3821	0.4	2	6066.1520	1.0	7				5085.2669	0.0	2
10	4	6	8297.0338	0.1	2				6186.8042	0.7	4				5210.4457	0.5	2
10	4	7	8289.1108	0.4	2				6180.9713	0.4	2				5206.2379	0.7	2
10	5	5	8403.7337	0.5	2				6317.0324		1						
10	5	6	8403.0557	0.6	2				6316.6074		1						
11	0	11	8116.1438		1	6214.7075	0.7	2	5966.4812	0.7	3	5025.3076		1	4973.4169		1
11	1	10	8243.6849		1	6339.2429	0.7	2	6107.7076	0.5	3						
11	1	11	8116.1768	1.0	2	6214.8009		1	5966.5590	0.4	3	5025.5859		1	4973.6084		1
11	2	9	8342.6533	1.0	4	6429.2249		1	6212.2719	1.0	4				5222.6605		1
11	2	10	8244.8156	1.0	3	6341.9657		1	6109.9697		1						
11	3	8	8408.2080	1.0	2				(222.1.10.1	1.0	2						
11	3	9	8355.2737	0.7	3				6233.1484	1.0	3						
11	4	7	84/3.4480	0.6	3				6363.3639	0.4	2						
11	4	8	8458.6422	0.0	1												
11	5	67	8572 8224	0.8	2												
11	5	12	0312.8234 8270 7276	20	1 ว	6369 0121		1	6121 0001	1.0	2	5181 1624		1	5127 0220		1
12	1	12	8603 1511	2.0	∠ 1	6505 8027		1 1	6277 9665	0.0	2	5101.4024		1	5127.7330		1
12	1	12	8270 7463		1	6368 0616		1	6121 9236	0.9	1	5181 6186		1	5128 0350		1
12	2	10	8525 2800		1	0500.0010		1	0121.7250		1	2101.0100		1	5120.0550		1
	-	10	0000.2000		*												

Table 7 (continued)

			(002)			(200)			(011)			(030)			(110)		
J	<i>K</i> <sub>a</sub> 1	$K_c$	E 2	$\frac{\delta}{3}$	N 4	E 5	δ 6	N 7	E 8	δ 9	N 10	Е 11	δ 12	N 13	E 14	δ 15	N 16
12	2	11	8411.8558		1				6279.2231	0.2	2						
12	3	9							6480.2985		1						
12	3	10	8533.1325		1				6413.6522		1						
12	4	8	8667.6481	1.0	2												
12	5	7	8761.8774		1												
12	5	8	8758.0914		1												
13	0	13													5294.5312		1
13	1	12				6684.3004		1									
13	1	13				6533.3922		1	6289.4987		1				5294.5741		1
13	2	11				6804.3172		1	6598.2236		2						
13	2	12	8591.1225		1				6460.8070	1.0	2						
13	3	10							6460.1230	1.0	1						
14	0	14	8616.4177		1				6469.2261		1				5473.1957		1
14	1	14	8616.4177		1												

Table 8

Spectroscopic parameters of the (000), (010), (020), (100), (001), (011), (200), and (002) vibrational states of HD<sup>18</sup>O (in cm<sup>-1</sup>)

	(000)	(010)	(020)	(100)	(001)	(011)	(200)	(002)
Ε	0	1396.26715(40)	2767.2030(21)	2709.2929(14)	3696.33022(36)	5071.49629(47)	5335.36189(98)	7229.18576(58)
A	23.110310(27)	25.16599(14)	27.63242(97)	23.02060(74)	22.08749(11)	23.97815(22)	22.6320(13)	21.09042(30)
В	9.059135(12)	9.192757(34)	9.30825(19)	8.89735(17)	9.049520(40)	9.189927(63)	8.70728(22)	9.039264(69
С	6.360901(12)	6.278102(26)	6.20879(20)	6.24851(10)	6.283315(31)	6.199903(40)	6.14933(15)	6.205176(61
$\Delta_K \times 10^2$	1.18744(38)	2.2309(12)	3.059(13)	2.534(12)	1.0845(11)	2.0654(26)	1.780(49)	0.9039(69)
$\Delta_{JK} \times 10^3$	1.2734(11)	1.1136(46)	0.489(23)	1.814(22)	1.1041(34)	0.9654(77)	2.238(26)	0.9322(67)
$\Delta_J \times 10^3$	0.35318(13)	0.40399(45)	0.4322(24)	3.774(15)	0.38582(40)	0.41599(69)	3.334(22)	0.3739(10)
$\delta_K \times 10^3$	2.0836(53)	3.4569(56)	4.390(62)	3.298(20)	1.9523(51)	3.329(12)	3.353(61)	1.8581(56)
$\delta_J \times 10^3$	0.118001(89)	0.14414(16)	0.15924(65)	0.13329(97)	0.12464(25)	0.15106(47)	0.1079(14)	0.13056(42)
$H_K \times 10^4$	0.4575(20)	1.2029(45)	0.070	6.996(77)	0.3858(40)	1.379(12)	2.42(67)	0.372(29)
$H_{KJ} \times 10^5$	-0.863(18)	-1.732(28)	-4.33(19)	0.142	-0.720(25)	-1.854(62)	_	-0.988(30)
$H_{JK} \times 10^6$	2.359(63)	4.707(74)	6.82(57)	5.46(17)	1.966(91)	5.90(20)	15.93(57)	2.833(90)
$H_J \times 10^6$	0.03708(83)	0.0556(15)	0.025	0.1154(73)	0.0437(15)	0.0407(39)	-0.0708	0.0260(94)
$h_K \times 10^5$	1.477(81)	4.447(62)	5.50(48)	6.87(29)	1.141(65)	4.63(24)	11.44(50)	1.603(96)
$h_{JK} \times 10^6$	1.069(39)	1.643(55)	1.08(35)	2.78(13)	0.906(52)	1.75(10)	1.82(33)	1.037(65)
$h_J \times 10^6$	1.658(50)	0.02677(53)	_	0.0658(39)	0.0208(12)	0.0224(23)	-0.00368	0.0198(18)
$L_K \times 10^6$	-0.174(11)	-0.6722(99)	-0.376(87)	-16.30(19)	-0.1609(78)	-1.354(42)	_	-0.210(75)
$L_{KKJ} \times 10$	$0^{7}$ 0.187	1.71(11)	_	_	0.791(95)	1.16(49)	19.1(17)	0.907
$L_{KJ} \times 10^7$	-0.0293	-0.514(44)	_	-2.41(33)	-0.324(34)	-0.419	-3.84(21)	-0.221(90)
$L_{KJJ} \times 10^{\circ}$	$^{9}-0.681$	-1.92(31)	_	_	2.25(41)	-4.00	_	-3.81(34)
$l_K \times 10^6$	-0.0324	-0.320(27)	_	-1.14(13)	-0.114(14)	-0.261(84)	_	-0.153(39)
$l_{KJ} \times 10^8$	-0.398	-1.39(26)	_	_	2.22(28)	-1.86(80)	_	_
$l_{JK} \times 10^9$	-0.134	-0.734	_	_	-0.720		_	_
$P_K \times 10^7$	0.00456(20)	0.0131(11)	_	-1.610(16)	0.00224(26)	0.0923(16)	-0.358	0.00696
$p_K \times 10^9$	0.0126	-0.549	_	14.0(12)	0.052		_	_
$\sigma \times 10^3$	0.194	1.06	4.97	3.32	0.924	1.09	1.82	1.33

## Acknowledgments

The authors are indebted to O. Naumenko of IAO (Russia) for suggestions and discussions. This work was jointly supported by the Natural Science Foundation of China (20473079 and 10574124), and by the Chinese Academy of Sciences.

# Appendix A. Supplementary data

Supplementary data for this article are available on ScienceDirect (www.sciencedirect.com) and as part of the

Ohio State University Molecular Spectroscopy Archives (http://msa.lib.ohio-state.edu/jmsa\_hp.htm).

#### References

- [1] G. Steenbeckeliers, J. Bellet, Compt. Rend. 273B (1971) 471–474.
- [2] F.X. Powell, D.R. Johnson, Phys. Rev. Lett. 24 (1970) 647.
- [3] F.C. De Lucia, P. Helminger, R.L. Cook, W. Gordy, Phys. Rev. Lett. A6 (1972) 1324–1326.
- [4] P. Helminger, F.C. De Lucia, J. Mol. Spectrosc. 70 (1978) 263–269.
- [5] F.C. De Lucia, P. Helminger, W.H. Kirchhoff, J. Phys. Chem. Ref. Data 3 (1974) 211–219.

- [6] F. Winther, J. Mol. Spectrosc. 65 (1977) 405-419.
- [7] R.A. Toth, J.-M. Flaud, C. Camy-Peyret, J. Mol. Spectrosc. 67 (1977) 206–218.
- [8] R.A. Toth, J.-M. Flaud, C. Camy-Peyret, J. Mol. Spectrosc. 67 (1977) 185–205.
- [9] P. Helminger, J. Mol. Spectrosc. 70 (1978) 263-269.
- [10] C. Camy-Peyret, J.-M. Flaud, N. Papineau, CR. Acad. Sci. B Phys. 290 (1980) 537–540.
- [11] C. Camy-Peyret, J.-M. Flaud, R.A. Toth, J. Mol. Spectrosc. 87 (1981) 233–241.
- [12] J.-P. Chevillard, J.-Y. Mandin, J.-M. Flaud, C. Camy-Peyret, Can. J. Phys. 63 (1985) 1112–1127.
- [13] A.D. Bykov, V.N. Saveliev, O.N. Ulenikov, J. Mol. Spectrosc. 118 (1986) 31–315.
- [14] J.-P. Chevillard, J.-Y. Mandin, C. Camy-Peyret, J.-M. Flaud, Can. J. Phys. 64 (1986) 746–761.
- [15] J.-P. Chevillard, J.-Y. Mandin, J.-M. Flaud, C. Camy-Peyret, Can. J. Phys. 65 (1987) 777–789.
- [16] O.N. Ulenikov, A.S. Zhilyakov, J. Mol. Spectrosc. 133 (1989) 239-243.
- [17] R.A. Toth, J. Opt. Soc. Am. B 9 (1992) 462-482.
- [18] R.A. Toth, J. Opt. Soc. Am. B 10 (1993) 1526-1544.
- [19] R.A. Toth, J. Mol. Spectrosc. 166 (1994) 184-203.
- [20] R.A. Toth, Appl. Opt. 33 (1994) 4869.
- [21] A. Bykov, O. Naumenko, T. Petrova, A. Scherbakov, L. Sinitsa, J.-Y. Mandin, C. Camy-Peyret, J.-M. Flaud, J. Mol. Spectrosc. 172 (1995) 243–253.
- [22] R.A. Toth, J. Mol. Spectrosc. 190 (1998) 379-396.
- [23] M. Tanaka, J.-W. Brault, J. Tennyson, J. Mol. Spectrosc. 216 (2002) 77–80.

- [24] M. Tanaka, O. Naumenko, J.W. Brault, J. Tennyson, J. Mol. Spectrosc. 234 (2005) 1–9.
- [25] M. Tanaka, M. Sneep, W. Ubachs, J. Tennyson, J. Mol. Spectrosc. 226 (2004) 1–6.
- [26] P. Macko, D. Romanini, S.N. Mikhailenko, O.V. Naumenko, S. Kassi, A. Jenouvrier, Vl.G. Tyuterev, A. Campargue, J. Mol. Spectrosc. 227 (2004) 90–108.
- [27] A.-W. Liu, O. Naumenko, K.-F. Song, B. Voronin, S.-M. Hu, J. Mol. Spectrosc. 236 (2006) 127–133.
- [28] A.-W. Liu, S.-M. Hu, C. Camy-Peyret, J.-Y. Mandin, O. Naumenko, B. Voronin, J. Mol. Spectrosc. 2006, in press.
- [29] L.S. Rothman, D. Jacquemart, A. Barbe, D. Chris Benner, M. Birk, L.R. Brown, M.R. Carleer, C. Chackerian Jr., K. Chance, L.H. Coudert, V. Dana, V.M. Devi, J.-M. Flaud, R.R. Gamache, A. Goldman, J.-M. Hartmann, K.W. Jucks, A.G. Maki, J.-Y. Mandin, S.T. Massie, J. Orphal, A. Perrin, C.P. Rinsland, M.A.H. Smith, J. Tennyson, R.N. Tolchenov, R.A. Toth, J. Vander Auwera, P. Varanasi, G. Wagner, J. Quant. Spectrosc. Radiat. Transf. 96 (2005) 139–204.
- [30] H. Partridge, D.W. Schwenke, J. Chem. Phys. 106 (1997) 4618-4639.
- [31] D.W. Schwenke, H. Partridge, J. Chem. Phys. 113 (2000) 6592-6597.
- [32] R.A. Roth, J. Mol. Spectrosc. 162 (1993) 20-40.
- [33] W.F. Wang, T.L. Tan, B.L. Tan, P.P. Ong, J. Mol. Spectrosc. 176 (1996) 226–228.
- [34] R.A. Toth, J. Mol. Struc. 742 (2005) 49–68.
- [35] J.K.G. Watson, J. Chem. Phys. 46 (1967) 1935-1949.
- [36] O.N. Ulenikov, G.A. Onopenko, I.M. Olekhnovitch, S. Alanko, V.-M. Horneman, M. Koivusaari, R. Anttila, J. Mol. Spectr. 189 (1998) 74–82.